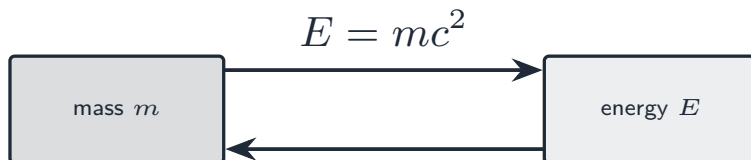


Nuclear physics

A-Level Physics

Mass-energy equivalence

Einstein's special relativity gives the famous link (**mass-energy equivalence** 质能等价):



Mass and energy are equivalent and can change into each other

$$E = mc^2,$$

where $c = 3.00 \times 10^8 \text{ m s}^{-1}$. A mass m matches an **energy** 能量 E — the two can change into each other. For a mass change Δm :

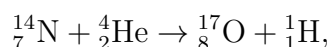
$$\Delta E = c^2 \Delta m.$$

In nuclear physics the masses are tiny but c^2 is huge, so a small mass change means a large energy. A mass change of 1 u ($1.661 \times 10^{-27} \text{ kg}$) matches $\Delta E \approx 1.49 \times 10^{-10} \text{ J}$. This gives a conversion you will use again and again:

$$1 \text{ u} \longleftrightarrow 931 \text{ MeV}.$$

Nuclear reactions

A **nuclear reaction** 核反应 is written like



with **nucleon number** 核子数 conserved (top numbers: $14 + 4 = 17 + 1$) and charge conserved (bottom numbers: $7 + 2 = 8 + 1$) — this is **conservation of charge** 电荷守恒. Use these to fill in an unknown: identify the species, then balance the top and bottom numbers.

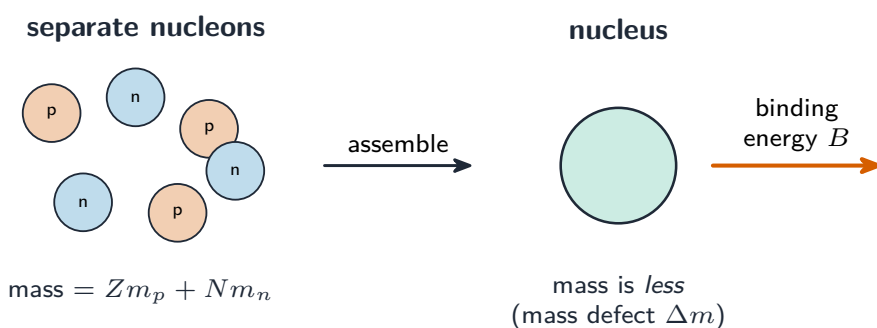
Mass defect and binding energy

The mass of a **nucleus** 原子核 is **less than** the total mass of its separate **protons** 质子 and **neutrons** 中子. The difference is the **mass defect** 质量亏损 Δm :

$$\Delta m = (Zm_p + Nm_n) - m_{\text{nucleus}}.$$

By $E = mc^2$, this "missing" mass was released as energy when the nucleus formed. To pull the nucleus fully apart you must put that energy back — the **binding energy** 结合能 B :

$$B = \Delta m \cdot c^2.$$



The assembled nucleus has less mass than its separate nucleons; the missing mass is released as binding energy

Worked example. A helium-4 nucleus has a mass defect of $\Delta m = 0.0304$ u. Find its binding energy. (1 u corresponds to 931 MeV.)

$$B = \Delta m \cdot c^2 = 0.0304 \times 931 \approx 28 \text{ MeV}.$$

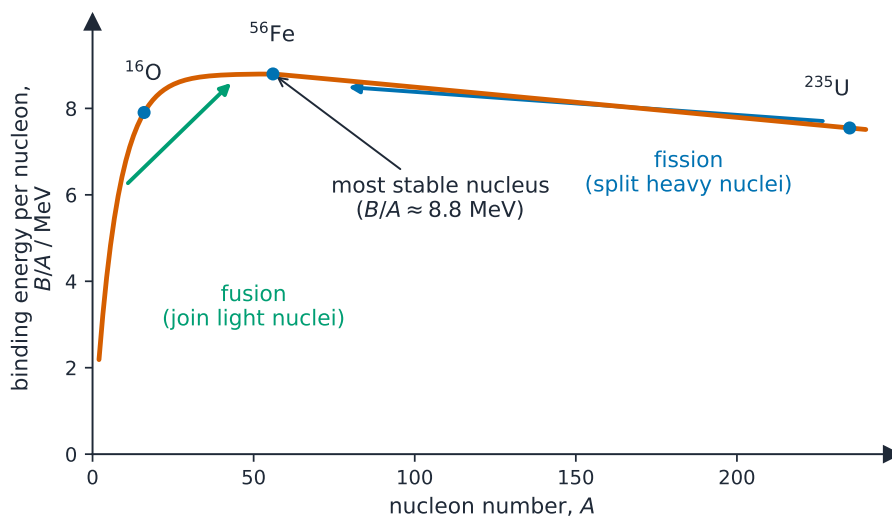
A more tightly bound nucleus has a larger mass defect and larger binding energy. The **binding energy per nucleon** 比结合能 is B/A (usually in MeV per nucleon) — a measure of how tightly each nucleon is held, useful for comparing nuclides.

Binding energy per nucleon vs nucleon number

A graph of B/A against A has a typical shape:

- for light nuclei ($A < 20$), B/A rises quickly (with a spike at the very stable ${}^4_2\text{He}$).
- around $A \sim 56$ (iron), B/A reaches its **maximum** of about 8.8 MeV. **Iron-56 is the most stable nucleus.**
- for heavy nuclei ($A > 100$), B/A falls slowly, to about 7.5 MeV for uranium.

So the curve is dome-shaped, rising to iron then falling.



binding energy per nucleon peaks near ${}^{56}\text{Fe}$

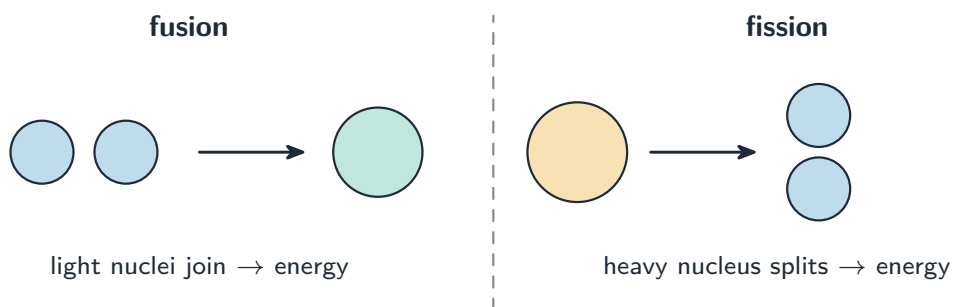
Binding energy per nucleon peaks near iron ($A \approx 56$); lighter and heavier nuclei are less tightly bound

Nuclear fusion and fission



A nuclear power station releases energy by nuclear fission.

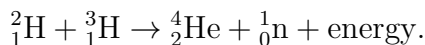
Energy is **released** when nuclei move **towards** the iron peak — by joining light nuclei or splitting heavy ones.



Fusion joins light nuclei; fission splits a heavy nucleus — both release energy by moving towards the iron peak

Nuclear fusion

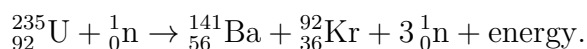
Nuclear fusion 核聚变 joins two light nuclei into one heavier nucleus:



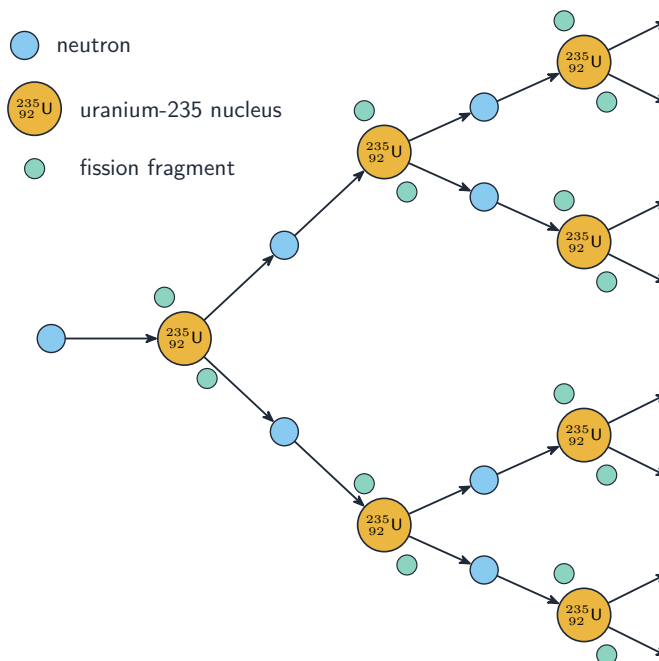
The product has greater binding energy per nucleon than the reactants, so energy is released. Fusion powers stars. It needs very high temperatures (millions of kelvin) so the nuclei have enough **kinetic energy** 动能 to beat their **electrostatic** 静电 repulsion and get close enough for the **strong nuclear force** 强核力 to take over.

Nuclear fission

Nuclear fission 核裂变 splits a heavy nucleus into two lighter ones:



The products have higher binding energy per nucleon than ${}^{235}_{92}\text{U}$, so energy is released. The extra neutrons can cause more fissions — a **chain reaction** 链式反应 in a large enough mass of fuel (the **critical mass** 临界质量). This is the basis of nuclear power and weapons.



In an uncontrolled chain reaction each fission of uranium-235 releases neutrons that cause more fissions

Calculating the energy released

1. find the total mass of the reactants.
2. find the total mass of the products.
3. mass change $\Delta m = m_{\text{reactants}} - m_{\text{products}}$ (positive when energy is released).
4. energy released $\Delta E = c^2 \Delta m$.

In kg this gives joules; in atomic mass units use $\Delta E \text{ (MeV)} = \Delta m \text{ (u)} \times 931$.

Worked example. In a nuclear reaction the total mass decreases by 0.020 u. Find the energy released.

$$\Delta E = \Delta m \text{ (u)} \times 931 = 0.020 \times 931 \approx 19 \text{ MeV.}$$

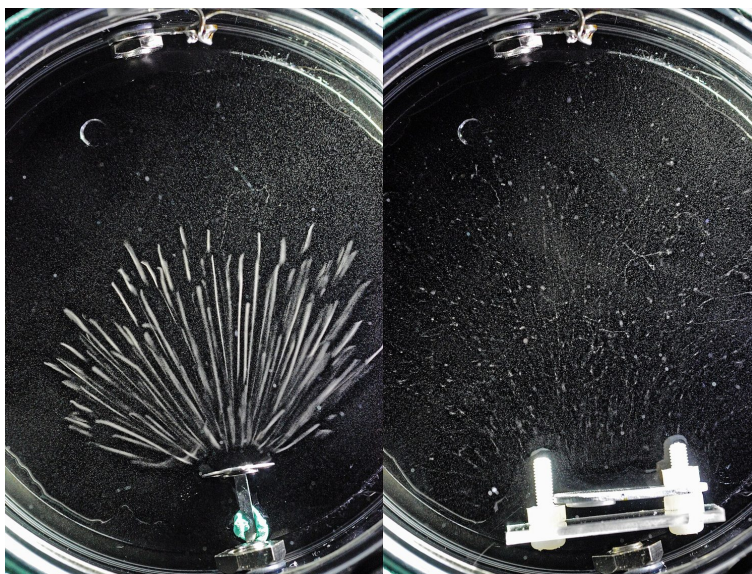
Radioactive decay

Random and spontaneous

Radioactive decay is:

- **spontaneous** 自发 — it happens with no outside trigger, and the rate is not changed by temperature, pressure or chemical state; and
- **random** 随机 — you cannot predict when a given nucleus will decay, only the probability that it decays in a time.

Evidence for randomness: the count rate **fluctuates**. A **Geiger counter** 盖革计数器 next to a source clicks at uneven intervals — never a steady stream — although the long-run mean rate is well-defined.



Each beta particle from the source leaves a thin track in a cloud chamber – direct evidence of separate, random decays

Activity and decay constant

For N undecayed nuclei of a **radionuclide** 放射性核素, the rate of decay is

$$A = \lambda N.$$

- A is the **activity** 活度 — decays per unit time. Unit: **becquerel** 贝克勒尔 (Bq) = s^{-1} .
- λ is the **decay constant** 衰变常数 — the probability per unit time that a nucleus decays. Unit: s^{-1} .

λ is fixed for a nuclide; a larger sample (larger N) has proportionally larger activity.

Exponential decay

Since λ is the fractional decay rate, $\frac{dN}{dt} = -\lambda N$, whose solution is an **exponential decay** 指数衰减:

$$N = N_0 e^{-\lambda t}.$$

Because $A = \lambda N$, the activity (and any **count rate** 计数率 proportional to it) decays the same way:

$$A = A_0 e^{-\lambda t}.$$

Why exponential? For each nucleus, λ is a fixed probability per unit time, independent of the others and of the nucleus's age. So the same **fraction** decays in each time interval, which gives exponential decay.

Half-life

The **half-life** 半衰期 $t_{1/2}$ is the time for the number of undecayed nuclei (or the activity, or the count rate) to fall to **half**. From $N = N_0 e^{-\lambda t}$ with $N = N_0/2$:

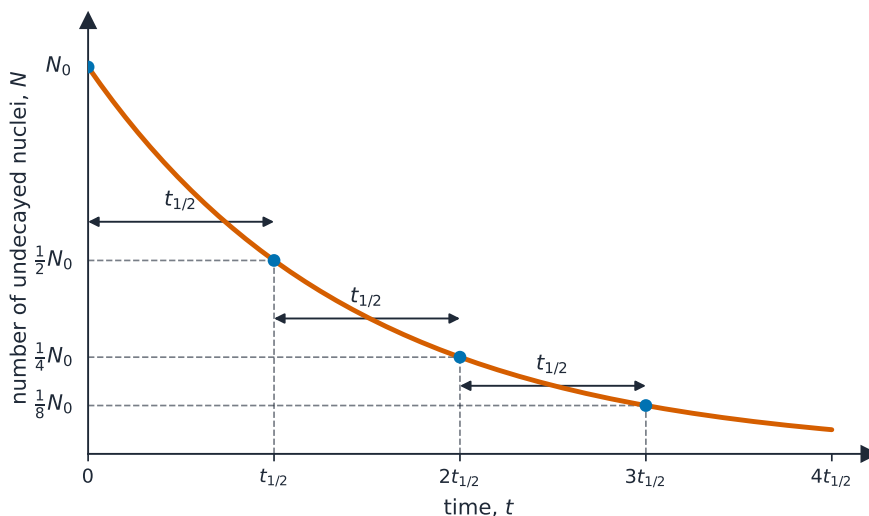
$$\ln 2 = \lambda t_{1/2}, \quad \lambda = \frac{\ln 2}{t_{1/2}} \approx \frac{0.693}{t_{1/2}}.$$

Worked example. A radioactive isotope has a half-life of 6.0 hours. Find its decay constant.

Convert the half-life to seconds: 6.0 h = 21 600 s. Then

$$\lambda = \frac{\ln 2}{t_{1/2}} = \frac{0.693}{21\,600} \approx 3.2 \times 10^{-5} \text{ s}^{-1}.$$

A larger decay constant means a shorter half-life. After n half-lives the surviving fraction is $(1/2)^n$; after 5 half-lives only about 3% remains.



$$N = N_0 e^{-\lambda t} \quad T_{1/2} = \ln 2 / \lambda$$

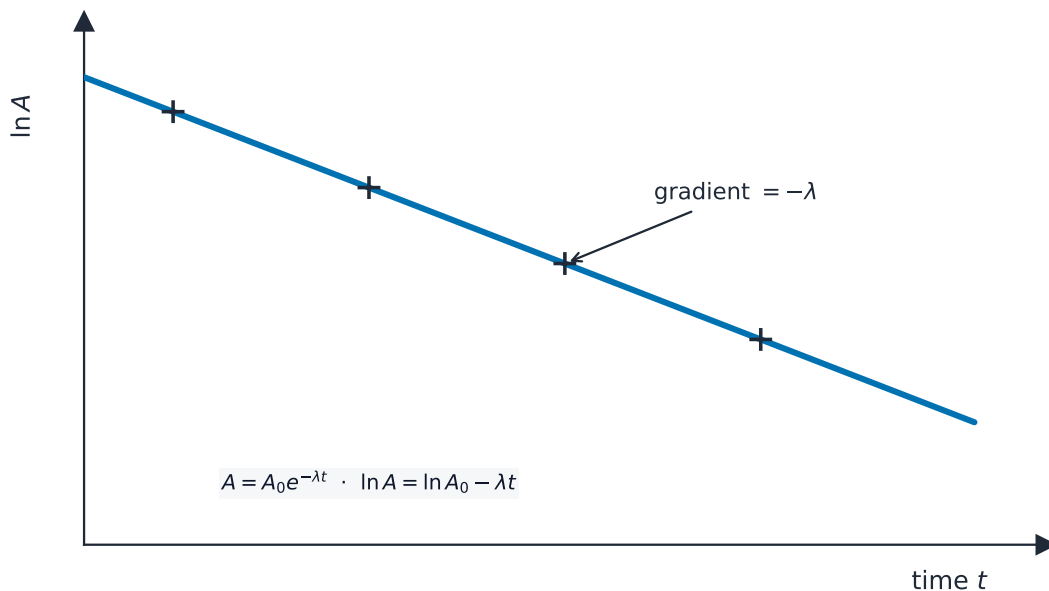
The number of undecayed nuclei falls by half in each half-life

Finding λ from data

Given A_0 and A at time t :

$$\lambda = \frac{1}{t} \ln \frac{A_0}{A}, \quad t_{1/2} = \frac{\ln 2}{\lambda}.$$

Taking logs of $A = A_0 e^{-\lambda t}$ gives $\ln A = \ln A_0 - \lambda t$, so a plot of $\ln A$ against t is a straight line with gradient $-\lambda$. Use this with several data points.



A graph of $\ln A$ against time is a straight line of gradient $-\lambda$

Exam tips

- $E = mc^2$; the **mass defect** is the mass lost when nucleons bind, released as binding energy.
- **Binding energy per nucleon** peaks near iron — both fusion (light nuclei) and fission (heavy nuclei) move towards it and release energy.
- Activity $A = \lambda N$; decay is exponential $N = N_0 e^{-\lambda t}$; half-life $t_{1/2} = \ln 2 / \lambda$.
- Balance nuclear equations by conserving **nucleon number and proton number**.